

Lecture 14. Physical meaning of the divergent series

$$\{ 1 + 2 + 3 + \dots = -\frac{1}{12} ? \text{ Why?}$$

Ramanujan's letter to Hardy

• { Casimir effect:

$$1^3 + 2^3 + 3^3 + \dots = \frac{1}{120}$$

① Ramanujan's letter to Hardy

$$S = 1 + 2 + 3 + 4 + 5 + 6$$

$$\begin{array}{r} 4S = 4 + 8 + 12 \\ \hline -3S = 1 - 2 + 3 - 4 + 5 - 6 = \frac{1}{(1+1)^2} = \frac{1}{4} \end{array}$$

$$S = -\frac{1}{12} !$$

$$\frac{1}{1+x} = 1 - x + x^2 - x^3$$

$$\begin{array}{r} \frac{1}{(1+x)^2} = 1 - x + x^2 - x^3 \\ x \quad 1 - x + x^2 - x^3 \\ \hline 1 - x + x^2 - x^3 \\ -x + x^2 - x^3 \\ \hline x^2 - x \\ -x^3 \end{array}$$

$$1 - 2x + 3x^2 - 4x^3$$

Can you believe it?

§3. Simplified version of Casimir effect

standing wave mode

$$\psi = \sin k_n x \quad \text{with} \quad k_n L = n\pi$$

$$\Rightarrow \text{Zero-point motion} \quad E_n = \hbar k_n c = \frac{\hbar c}{L} \pi \left(n + \frac{1}{2} \right)$$

↖ zero point motion

$$\text{Hence} \quad E_0(L) = \frac{1}{2} \cdot 2 \sum_{n=1}^{\infty} \hbar \omega_n = \frac{\hbar c \pi}{L} (1 + 2 + 3 + \dots)$$

obviously, this summation diverges. Does it really have any significance?

of course, the absolute value of $E_0(L)$ does not have meaning, we care about the energy changes as $L \rightarrow L + \delta L$. But even

though $F = -\frac{\partial}{\partial L} E(L) = \frac{\hbar c \pi}{L^2} \sum_{n=1}^{\infty} n$ ← this still does not make sense!

Even the sign is wrong! One would expect the force is repulsive, but as we will see, the force is attractive!

The reason is. that we cannot actually sum over to infinity. In a ~~metallic~~ metallic plate, there exist a cut off frequency ω_p (typically the plasmon frequency). Above ω_c , the

E & M wave would penetrate the plate, i.e. the plate is

(4)

Hence, we should add the cut off explicitly in the summation:

$$\frac{\hbar c \pi}{L} n < \hbar \omega_p, \text{ i.e. } n < N(L) = \frac{\omega_p}{c} \frac{L}{\pi} = \frac{L}{\ell}, \text{ where } \ell = \frac{c \pi}{\omega_p}.$$

ℓ is a microscopic length scale.

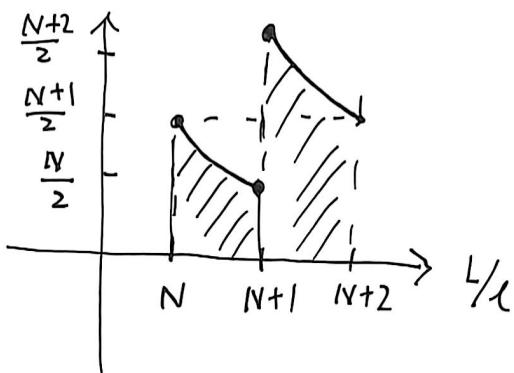
We define $S(\frac{L}{\ell}) = \left(\sum_{n=1}^{\lfloor \frac{L}{\ell} \rfloor} n \right) \frac{1}{\lfloor \frac{L}{\ell} \rfloor}$ $\lfloor \frac{L}{\ell} \rfloor$ is the integer $\leq \frac{L}{\ell}$.

$$\Rightarrow S(\frac{L}{\ell}) = \frac{1}{L} \cdot \frac{N(N+1)}{2} = \frac{\ell}{L} \cdot \frac{1}{2} \lfloor \frac{L}{\ell} \rfloor (\lfloor \frac{L}{\ell} \rfloor + 1).$$

Consider $N+1 > \frac{L}{\ell} > N \Rightarrow S(\frac{L}{\ell}) = \frac{1}{2} N(N+1)/(\lfloor \frac{L}{\ell} \rfloor)$

$$N+2 > \frac{L}{\ell} > N+1 \quad \frac{1}{2} (N+1)(N+2)/(\lfloor \frac{L}{\ell} \rfloor)$$

$S(\frac{L}{\ell})$

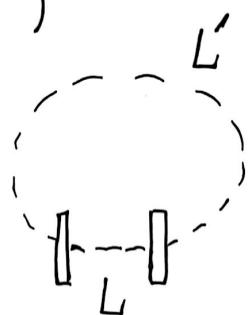


$S(\frac{L}{\ell})$ needs to be smoothed.

$$\begin{aligned} \tilde{S}(\frac{L}{\ell}) &= \int_N^{N+1} d(\frac{L}{\ell}) \cdot \frac{1}{\lfloor \frac{L}{\ell} \rfloor} \cdot \frac{N(N+1)}{2} = \frac{N(N+1)}{2} \ln \frac{N+1}{N} \\ &= \frac{N(N+1)}{2} \left(\frac{1}{N} - \frac{1}{2N^2} + \frac{1}{3N^3} + \dots \right) \\ &= \frac{N+1}{2} - \frac{N+1}{4N} + \frac{N+1}{6N^2} + \dots = \frac{N}{2} + \frac{1}{4} - \frac{1}{12N} + \dots \\ &= \frac{1}{2} \frac{L}{\ell} + \frac{1}{4} - \frac{1}{12} \frac{\ell}{L} \end{aligned}$$

$$E(L) = \frac{\hbar c \pi}{\ell} \cdot \underbrace{\frac{1}{4\ell} \sum_{n=1}^{[4\ell]} n}_{\text{Smoothed}} = \frac{\hbar c \pi}{\ell} \tilde{S}(4\ell)$$

$$= \frac{\hbar c \pi}{\ell} \left(\frac{1}{2} \frac{L}{\ell} + \frac{1}{4} - \frac{1}{12} \frac{1}{L} + \dots \right)$$



The EM wave also exists outside

the plates. We think it a 1D ring

$$E(L') = \frac{\hbar c \pi}{\ell} \left(\frac{1}{2} \frac{L'}{\ell} + \frac{1}{4} - \frac{1}{12} \frac{1}{L'} + \dots \right) \quad \text{with } L+L' = L_{\text{tot}}$$

$$\Rightarrow E_{\text{tot}} = \frac{\hbar c \pi}{\ell} \left(\frac{1}{2} \frac{L_{\text{tot}}}{\ell} + \frac{1}{2} - \frac{1}{12} \left(\frac{1}{L} + \frac{1}{L'} \right) + \dots \right) \quad \begin{matrix} \leftarrow^{\text{cancel}} \\ \text{vanish} \end{matrix}$$

Then
$$\boxed{F = - \frac{\partial E_{\text{tot}}}{\partial L} = - \frac{\hbar c \pi}{12} \frac{\partial}{\partial L} \left(\frac{1}{L} \right) = - \frac{\hbar c \pi}{12} \frac{1}{L^2}}$$

* A more sophisticated regularization method

We introduce the convergence factor $t = \frac{1}{N(L)} = \frac{\cancel{\hbar c \pi}}{L} l$

$$S = \sum_{n=1}^{\infty} n e^{-tn} \quad \leftarrow \text{effectively}$$

$$= - \frac{d}{dt} \sum_{n=1}^{\infty} e^{-tn} = - \frac{d}{dt} \frac{1}{1 - e^{-t}}$$

$$\frac{1}{1-e^{-t}} = \left(t - \frac{t^2}{2} + \frac{t^3}{3!} - \frac{t^4}{4!} \right)^{-1} = t^{-1} \underbrace{\left(1 - \frac{t}{2} + \frac{t^2}{6} - \frac{t^3}{24} + \frac{t^4}{120} \right)^{-1}}$$

$$\begin{aligned} & 1 + \frac{t}{2} - \frac{t^2}{6} + \frac{t^3}{24} - \frac{t^4}{120} \\ & \frac{t^2}{4} - \frac{t^3}{6} + \left(\frac{1}{36} + \frac{1}{24} \right) t^4 \\ & + \frac{t^3}{8} \quad \text{use } \frac{1}{2^2} \cdot \frac{1}{6} t^4 \\ & \hline 1 + \frac{t}{2} + \frac{t^2}{12} + 0 - \frac{1}{720} t^4 \end{aligned}$$

$$S = - \frac{d}{dt} \left(\frac{1}{t} + \frac{1}{2} + \frac{t}{12} - \frac{t^3}{720} + \dots \right) \Big|_{t \rightarrow 0}$$

$$= \frac{1}{t^2} - \frac{1}{12} + O(t^2)$$

$$\Rightarrow E_0(L) = \frac{\hbar C \pi}{L} \left(\dots - \left(\frac{L}{\zeta} \right)^2 - \frac{1}{12} + \dots \right)$$

$$= \frac{\hbar C \pi}{L^2} L - \underbrace{\frac{1}{12} \frac{\hbar C \pi}{L}}$$

but the

\uparrow this part is the same

UV part \approx

depends on regularization

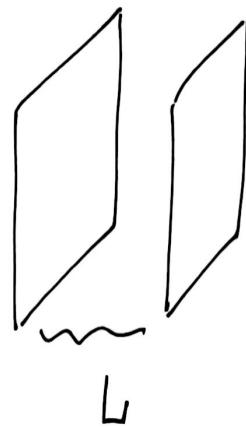
$$\text{Similarly, we have } F = - \frac{\hbar C \pi}{12} \frac{\partial}{\partial L} \left(-\frac{1}{L} \right) = - \frac{\hbar C \pi}{12} \frac{1}{L^2}$$

① 3D Casimir force

(7)

$$\psi_n(x, y, z, t) = \bar{e}^{-iw_n t} e^{i(k_x x + k_y y)} \sin k_z z$$

$$\omega_n = c \sqrt{k_x^2 + k_y^2 + \frac{n^2 \pi^2}{L^2}}$$



$$E = 2 \cdot \frac{\hbar}{2} A \int \frac{dk_x dk_y}{(2\pi)^2} \sum_{n=1}^{\infty} \omega_n$$

$$\frac{\omega_p^2}{c^2} - \frac{n^2 \pi^2}{L^2}$$

$$\frac{E}{A} = \frac{\hbar c}{A} \int \frac{dk_x dk_y}{(2\pi)^2} \underbrace{\sum_{n=1}^{\infty} (k_x^2 + k_y^2 + \frac{n^2 \pi^2}{L^2})}_{< \frac{\omega_p^2}{c^2}} = \frac{\hbar c}{4\pi} \int_0^{\omega_p} dq^2 \sum_{n=1}^{\infty} (q^2 + \frac{n^2 \pi^2}{L^2})^{1/2}$$

$$k_x^2 + k_y^2 + \frac{n^2 \pi^2}{L^2} < \frac{\omega_p^2}{c^2}$$

$$\sum_{n=1}^{\infty} \int_0^{\omega_p} dy \left(y + \frac{n^2 \pi^2}{L^2} \right)^{1/2}$$

$$= \frac{2}{3} \left[\frac{\omega_p^3}{c^3} - \left(\frac{n\pi}{L} \right)^3 \right]$$

$$\Rightarrow \frac{E}{A} = \text{const} - \frac{2}{3} \frac{\hbar c \pi^3}{4\pi L^3} \sum_{n=1}^{\infty} n^3$$

$$= - \frac{\hbar c \pi^2}{6 L^3} \sum_{n=1}^{\infty} n^3 + \text{const}$$

$$n < \frac{\omega_p L}{c \pi} =$$

~~$$S(\frac{l}{\ell}) = \left[\sum_{n=1}^{[\frac{l}{\ell}]} n^3 \right] = (\frac{l}{\ell})^3$$~~

~~$$= \left(\frac{l}{\ell} \right)^3 \frac{N^2(N+1)^2}{4} = \left(\frac{l}{\ell} \right)^3 \cdot \frac{1}{4} [\frac{l}{\ell}]^2 [\frac{l}{\ell} + 1]^2$$~~

Define

$$S = \sum_{n=1}^{\infty} n^3 e^{-tn} = -\frac{d^3}{dt^3} \sum_{n=1}^{\infty} e^{-tn}$$

$$= -\frac{d^3}{dt^3} \left(\frac{1}{t} + \frac{1}{2} + \frac{t}{12} - \frac{t^3}{720} + \dots \right) \Big|_{t \rightarrow 0}.$$

$$= \frac{6}{t^4} + \frac{1}{120} + \dots$$

$$\Rightarrow \frac{E_0(L)}{A} = -\frac{\hbar c}{6} \frac{\pi^2}{L^3} \left[\left(\frac{L}{\ell} \right)^4 + \frac{1}{120} \right]$$

$$= -\frac{\hbar c}{6} \frac{\pi^2}{\ell^4} L^4 - \frac{\hbar c}{720} \frac{\pi^2}{L^3}$$

$$\Rightarrow \frac{F}{A} = -\frac{\partial E}{\partial L} = -\frac{\hbar c \pi^2}{720} \frac{\partial}{\partial L} \left(\frac{1}{L^3} \right)$$

$$\frac{F}{A} = -\frac{\hbar c \pi^2}{240} \frac{1}{L^4}$$