

Lecture 3 Dimensional analysis (III)

§1 Unit systems (heat, temperature)

§2. E & M — SI units

— Gaussian units.

§3. natural units.

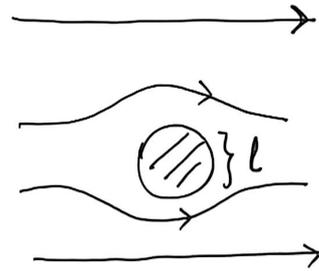
§ unit systems

Rayleigh's problem: Heat conduction in a flow.

Incompressible viscosity free fluid.

The object's size l , the temperature difference relative to the infinity $\Delta\theta$,

fluid velocity v , heat capacity per volume C , heat conductivity κ . We need to analyze heat releasing rate \hat{H} from the object to the fluid.



① Rayleigh's method.

He treats temperature and heat as fundamental quantities, then there're five fundamental ones, (M, L, T, Θ , Q).

Then $[l] = L$, $[v] = LT^{-1}$, $[\theta] = \Theta$, $[C] = Q\Theta^{-1}L^{-3}$, $[\kappa] = QL^{-1}\Theta^{-1}T^{-1}$

$$\frac{Q}{A \cdot t} = \kappa \cdot \nabla T$$

$$[H] = QT^{-1}$$

Then	M	L	T	Θ	Q
l	0	1	0	0	0
v	0	1	-1	0	0
$\Delta\theta$	0	0	0	1	0
C	0	-3	0	-1	1
κ	0	-1	-1	-1	1
H	0	0	-1	0	1

The first column is all 0's.

Hence, the rank is 4. There're $6 - 4 = 2$ dimensionless quantities. The first one

$$\left. \begin{aligned} \pi_1 &= \frac{H}{\kappa l \Delta\theta} \\ \pi_2 &= \frac{lvC}{\kappa} \end{aligned} \right\} \Rightarrow H = \kappa l \Delta\theta \Phi\left(\frac{lvC}{\kappa}\right)$$

Nevertheless, now we know heat and energy are the same, we should be able to use the same unit. Then we have M, L, T, Θ 4 fundamental units.

$$[L] = L, [v] = LT^{-1}, [\theta] = \Theta, [C] = [M][L]^{-1}[T]^{-2}[\Theta]^{-1}$$

$$[K] = [M][L][T]^{-3}[\Theta]^{-1}, [H] = [M][L]^2[T]^{-3}$$

	M	L	T	Θ
l	0	1	0	0
v	0	1	-1	0
$\Delta\theta$	0	0	0	1
C	1	-1	-2	-1
K	1	1	-3	-1
H	1	2	-3	0

\Rightarrow

still

$$\pi_1 = \frac{H}{K \Delta\theta}$$

$$\pi_2 = \frac{lvC}{K}, \text{ which is the same as before.}$$

Now Q is absent, but M appears.

Whether we should treat Q as an independent quantity, depends on our viewpoint. If it's viewed as independent, we should introduce the coefficient of mechanical equivalent of heat $J = 4.2 \text{ J/cal}$.

Then we add it into table $[J] = [M][L]^2[T]^{-2}[\Theta]^{-1}$

	M	L	T	Θ	Q
J	1	2	-2	0	-1
l	0	1	0	0	0
v	0	1	-1	0	0
$\Delta\theta$	0	0	0	1	0
C	0	-1	-2	-1	0
K	0	-1	-3	-1	0
H	0	0	-3	0	0

\rightarrow

	M	L	T	Θ	Q
l	0	1	0	0	0
v	0	1	-1	0	0
$\Delta\theta$	0	0	0	1	0
$C \cdot J$	1	-1	-2	-1	0
$K \cdot J$	1	1	-3	-1	0
$H \cdot J$	1	2	-3	0	0

if we absorb J 's unit,

The next step question is whether we should treat ~~temperature~~ and energy as the same. If they're different, we should have $k_B T = E$.

$$[k_B] = [M][L][T]^{-2} [\Theta]^{-1} \leftarrow \text{trade off replace } [\Theta] \text{ by } [M][L^2][T^{-2}]$$

	M	L	T	Θ	Q		M	L	T	Θ	Q	
k_B	1	2	-2	-1	0							
J	1	2	-2	0	-1	→ transformation	l	0	1	0	0	
C	0	1	0	0	0		v	0	1	-1	0	0
$\Delta\theta$	0	0	0	1	0		$k_B \Delta\theta$	1	2	-2	0	0
C	0	-3	0	-1	1		J C/k _B	0	-3	0	0	0
K	0	-1	-1	-1	1		J K/k _B	0	-1	-1	0	0
H	0	0	-1	0	1		J H	1	2	-3	0	0

Now we have $6-3=3$ independent unit. Hence in principle,

Rayleigh was not so correct. We can have

$$\pi_1 = \frac{JH}{\frac{JK}{k_B} l \cdot k_B \Delta\theta} = \frac{H}{k l \Delta\theta}, \quad \pi_2 = \frac{l v (J C/k_B)}{(J K/k_B)} = \frac{l v c}{k}$$

$$\pi_3 = \frac{J C}{k_B} l^3$$

Hence, we arrive at $H = k l \Delta\theta \Phi\left(\frac{l v c}{k}, \frac{J C}{k_B} l^3\right)$.

It seems more reasonable. Nevertheless, we are not considering stat-mech, and no heat and mechanical energy transfer. Hence

k_B and J should not appear. We recover $H = k l \Delta\theta \Phi\left(\frac{l v c}{k}\right)$.

{ E & M units

Whether we should introduce a new unit for E & M? In the SI units, electric current I is defined as a fundamental unit, $[I][T] = [\text{charge}]$, then it's linearly independent from M, L and T. For convenience, we use charge in stead of current. If we use charge, denoted as 'Q' (not heat in this section!)

Then $F = \frac{1}{4\pi\epsilon_0} \frac{e^2}{r^2}$, An additional constant ϵ_0 is

introduced, $[F][\epsilon_0] = [e^2][L^{-2}]$ $[F] = [M][L][T]^{-2}$

$$\Rightarrow [\epsilon_0] = Q^2 M^{-1} L^{-3} T^2$$

electric field $[E] = [F]/[e] = Q^{-1} M L T^{-2}$

potential voltage $[U] = [E]L = Q^{-1} M L^2 T^{-2}$

electric current $[I] = Q T^{-1}$

density $[j] = Q L^{-2} T^{-1}$

conductivity $[\sigma] = [j]/[E] = Q^2 M^{-1} L^{-3} T$

Magnetic field $\vec{F} = e\vec{v} \times \vec{B} \Rightarrow [B] = \frac{[F][v]^{-1}}{[Q]^{-1}} = Q^{-1} M T^{-1}$

Ampère's law $\oint B \cdot dl = \mu_0 I \Rightarrow [\mu] = [B] \cdot L / Q \cdot T = Q^{-2} M L$

→ energy density

$$U = \frac{1}{2} \epsilon_0 E^2 + \frac{1}{2\mu_0} B^2$$

$$[U] = M L^{-1} T^{-2}$$

$$c^2 = \frac{1}{\mu_0 \epsilon_0}$$

Nevertheless, ϵ_0, μ_0 are too cumbersome. Physicists do not like them -

E and M are not symmetric either. Rather, we absorb $\frac{e^2}{4\pi\epsilon_0} \rightarrow e'^2$

and define $\vec{F} = q(\vec{E} + \frac{\vec{v}}{c} \times \vec{B})$, such that E and B carry the same unit.

The correspondence is as follows: $F = \frac{e^2}{4\pi\epsilon_0 r^2} \Leftrightarrow F' = \frac{e'^2}{r^2} = F$

$$\Rightarrow e' = \frac{e}{\sqrt{4\pi\epsilon_0}} \quad E = \frac{e}{4\pi\epsilon_0 r^2} \quad E' = \frac{e'}{r^2} = \sqrt{4\pi\epsilon_0} E$$

(those with a prime, means it's in Gaussian unit).

$$e\vec{v} \times \vec{B} = e'\frac{\vec{v}}{c} \times \vec{B}'$$

$$\Rightarrow B' = c \frac{e}{e'} B = c\sqrt{4\pi\epsilon_0} B.$$

Hence, the formula in the SI

$$= \sqrt{\frac{4\pi}{\mu_0}} B$$

unit

$$\mathcal{U} = \frac{1}{2} \epsilon_0 E^2 + \frac{1}{2\mu_0} B^2 = \frac{1}{8\pi} (E'^2 + B'^2)$$

$$\left\{ \begin{array}{l} E = \frac{1}{\sqrt{4\pi\epsilon_0}} E' \\ B = \sqrt{\frac{\mu_0}{4\pi}} B' \end{array} \right.$$

This set of unit is called the Gaussian unit.

$$j = \frac{e}{\Delta t A} \quad j' = \frac{e'}{\Delta t A} = \frac{1}{\sqrt{4\pi\epsilon_0}} j$$

Now check Maxwell equation

$$\nabla \cdot \vec{E} = \rho/\epsilon_0 \rightarrow \nabla \cdot \frac{1}{\sqrt{4\pi\epsilon_0}} \vec{E}' = \sqrt{4\pi\epsilon_0} \rho'/\epsilon_0 \Rightarrow \nabla \cdot \vec{E}' = 4\pi \rho'$$

$$\nabla \cdot \vec{B} = 0 \rightarrow \nabla \cdot \vec{B}' = 0$$

$$\nabla \times \vec{E} = -\frac{\partial \vec{B}}{\partial t} \rightarrow \nabla \times \frac{1}{\sqrt{4\pi\epsilon_0}} \vec{E}' = -\frac{\partial}{\partial t} \sqrt{\frac{\mu_0}{4\pi}} \vec{B}' \Rightarrow \nabla \times \vec{E}' = -\frac{1}{c} \frac{\partial \vec{B}'}{\partial t}$$

$$\nabla \times \vec{B} = \mu_0 \vec{j} + \frac{1}{c^2} \frac{\partial \vec{E}}{\partial t} \rightarrow \nabla \times \sqrt{\frac{\mu_0}{4\pi}} \vec{B}' = \mu_0 \sqrt{4\pi\epsilon_0} j' + \frac{1}{c^2} \frac{\partial}{\partial t} \frac{1}{\sqrt{4\pi\epsilon_0}} \vec{E}'$$

$$\nabla \times \vec{B}' = 4\pi \sqrt{\mu_0 \epsilon_0} j' + \frac{1}{c^2} \frac{1}{\sqrt{\mu_0 \epsilon_0}} \frac{\partial \vec{E}'}{\partial t} \Rightarrow \nabla \times \vec{B}' = 4\pi j' + \frac{1}{c^2} \frac{\partial \vec{E}'}{\partial t}$$

{ Natural unit (SI)

atomic length $a = \frac{\hbar^2}{m e^2} \cdot 4\pi\epsilon_0 \simeq 0.5 \text{ \AA}, \quad \frac{\hbar^2}{m a^2} = \frac{e^2}{4\pi\epsilon_0 a}$

Rydberg energy $E = \frac{e^2}{4\pi\epsilon_0 a} = \frac{m e^4}{(4\pi\epsilon_0)^2 \hbar^2} = 27.2 \text{ V}$

	Q	M	L	T
$4\pi\epsilon_0$	2	-1	-3	2
e^2	2	0	0	0
\hbar	0	1	2	-1
c	0	0	1	-1
m	0	1	0	0
I	1	0	0	-1
u	-1	1	2	-2
R	-2	1	2	-1
C	2	-1	-2	2
ϕ	-1	1	2	-1
L	-2	1	2	0
E	-1	1	1	-2
B	-1	1	0	-1
$\frac{\mu_0}{4\pi}$	-2	1	1	0
P	1	0	1	0
磁矩 $\mu = IA$	1	0	2	-1
$\frac{\mu_0}{4\pi}$	-2	1	1	0

$\Rightarrow \pi_1 = \frac{e^2}{4\pi\epsilon_0 \hbar c} = \frac{1}{137} = \alpha$

$\alpha a = \frac{\hbar}{m c} \quad \therefore \text{Compton length}$

$\alpha^2 a = \frac{1}{4\pi\epsilon_0} \frac{e^2}{m c} \quad \therefore \text{electron classic radius}$

Quantum unit of R, C, L

$R: R_Q = \hbar/e^2 = \frac{6.6 \times 10^{-34}}{1.6^2 \times 10^{-38}} = 25.8 \text{ k}\Omega$

$C: \left[\frac{C}{4\pi\epsilon_0} \right] = L \Rightarrow \frac{C}{4\pi\epsilon_0} = \frac{\hbar^2}{m e^2} \cdot 4\pi\epsilon_0$

$\Rightarrow \frac{C_Q}{a} = a \cdot 4\pi\epsilon_0 = \frac{0.5 \times 10^{-10}}{9 \times 10^9} \text{ F}$
 $= 5.5 \times 10^{-24} \text{ F}$

$L: [L \cdot 4\pi\epsilon_0] = L^{-1} T^2 \Rightarrow = L \left(\frac{T}{L} \right)^2$

$L_Q \cdot 4\pi\epsilon_0 = \frac{\hbar^2}{m e^2} \cdot 4\pi\epsilon_0 \frac{1}{c^2}$

$L_Q = \frac{\hbar^2}{m e^2} \frac{1}{c^2} = \frac{a}{4\pi\epsilon_0} \frac{1}{c^2} = L$

$= \frac{0.5 \times 10^{-10} \times 9 \times 10^9}{9 \times 10^{16}} \text{ H} = 5 \times 10^{-18} \text{ H}$

$\omega_Q = \frac{1}{\sqrt{L_Q C_Q}} = \frac{C}{a} \simeq 6 \times 10^{18} \text{ Hz}$

Gaussian unit Q is not an independent unit

	Q	M	L	T
$4\pi\epsilon_0$	2	-1	-3	2
$\mu_0/4\pi$	-2	1	1	0
$e^2 = e^2/4\pi\epsilon_0$	0	1	3	-2
\hbar	0	1	2	-1
c	0	0	1	-1
m	0	1	0	0
$I' = \frac{I}{\sqrt{4\pi\epsilon_0}}$	0	1/2	3/2	-2
$U' = U\sqrt{4\pi\epsilon_0}$	0	1/2	1/2	-1
$\phi' = \frac{\phi}{\sqrt{\frac{\mu_0}{4\pi}}}$	0	1/2	3/2	-1
$E' = E\sqrt{4\pi\epsilon_0}$	0	1/2	-1/2	-1
$B' = \frac{B}{\sqrt{\frac{\mu_0}{4\pi}}}$	0	1/2	-1/2	-1
$R' = R4\pi\epsilon_0$	0	0	-1	1
$C' = \frac{C}{4\pi\epsilon_0}$	0	0	1	0
$L' = L4\pi\epsilon_0$	0	0	-1	2

	Q	M	L	T
$P' = \frac{P}{\sqrt{4\pi\epsilon_0}}$	0	1/2	5/2	-1
$m' = m\sqrt{\frac{\mu_0}{4\pi}}$	0	1/2	5/2	-1

$$\begin{aligned}
 L'I' &= LI\sqrt{4\pi\epsilon_0} \\
 &= \phi\sqrt{4\pi\epsilon_0} \\
 &= \phi'\sqrt{\frac{\mu_0}{4\pi}} \cdot 4\pi\epsilon_0 = \frac{\phi'}{c}
 \end{aligned}$$

$$\Rightarrow L'I' = \frac{\phi'}{c} \rightarrow \text{Gaussian unit.}$$

$$\Rightarrow a = \frac{\hbar^2}{me^2}, \quad E = \frac{me^2}{\hbar^2}, \quad \alpha = \frac{e^2}{\hbar c}, \quad \alpha a = \frac{\hbar}{mc}, \quad \alpha^2 a = \frac{e^2}{mc}$$

$$R'_a = \hbar/e^2, \quad C'_a = a, \quad L'_a = L \cdot 4\pi\epsilon_0 = \frac{a}{c^2}$$

It's much simpler