

Lecture 8 Curvature tensor and Einstein's Equation

{ Geodesic deviation

{ Curvature tensor and symmetries

$$R_{\mu\nu\lambda\sigma} = -R_{\nu\mu\lambda\sigma} = -R_{\mu\nu\sigma\lambda} = R_{\lambda\sigma\mu\nu}$$

$$R_{\mu\nu\lambda\sigma} + R_{\mu\lambda\nu\sigma} + R_{\mu\sigma\nu\lambda} = 0$$

{ Ricci tensor

$$R_{\mu\nu} = g^{\lambda\sigma} R_{\mu\lambda\nu\sigma}, \quad R = g^{\mu\nu} R_{\mu\nu}$$

{ Construction of Einstein Equation

~~$$D_\mu G^{\mu\nu} = 0$$~~

$$G^{\mu\nu} = R^{\mu\nu} - \frac{1}{2} g^{\mu\nu} R = 8\pi G_N T^{\mu\nu}$$

§ Geodesic deviation

Consider a group of geodesics $\chi^M(z, h)$ where z is the proper time along a geodesic, and h distinguishes different geodesics.

$$\frac{\partial^2 \chi^M(z, h)}{\partial z^2} + \Gamma_{\nu\lambda}^M \frac{\partial \chi^\nu}{\partial z} \frac{\partial \chi^\lambda}{\partial z} = 0$$

Define tangent vector $u^M(z, h) = \partial \chi^M(z, h) / \partial z$, then

the geodesic becomes $\partial^2 u^M(z, h) / \partial z^2 + \Gamma_{\nu\lambda}^M u^\nu u^\lambda = 0.$

Define $v^M(z, h) = \partial \chi^M(z, h) / \partial h \Rightarrow$

$$\begin{aligned} \frac{\partial v^M}{\partial z} &= \frac{\partial^2 \chi^M}{\partial z \partial h} = \frac{\partial u^M}{\partial h} \Rightarrow \frac{Dv^M}{Dz} = \frac{\partial v^M}{\partial z} + \Gamma_{\nu\lambda}^M \frac{\partial \chi^\nu}{\partial z} v^\lambda \\ &= \frac{\partial u^M}{\partial h} + \Gamma_{\nu\lambda}^M u^\nu v^\lambda \end{aligned}$$

$$\begin{aligned} \frac{D^2 v^M}{Dz^2} &= \frac{\partial}{\partial z} \left(\frac{Dv^M}{Dz} \right) + \Gamma_{\rho\sigma}^M \frac{\partial \chi^\rho}{\partial z} \frac{Dv^\sigma}{Dz} \\ &= \frac{\partial}{\partial z} \left(\frac{\partial u^M}{\partial h} + \Gamma_{\nu\lambda}^M u^\nu v^\lambda \right) + \Gamma_{\rho\sigma}^M \frac{\partial \chi^\rho}{\partial z} \left(\frac{\partial u^\sigma}{\partial h} + \Gamma_{\nu\lambda}^\sigma u^\nu v^\lambda \right) \\ &= \frac{\partial}{\partial z} \frac{\partial u^M}{\partial h} + \frac{\partial}{\partial z} \Gamma_{\nu\lambda}^M u^\nu v^\lambda + \Gamma_{\nu\lambda}^M \frac{\partial u^\nu}{\partial z} v^\lambda + \Gamma_{\nu\lambda}^M u^\nu \\ &\quad + \Gamma_{\rho\sigma}^M \frac{\partial \chi^\rho}{\partial z} \frac{\partial u^\sigma}{\partial h} + \Gamma_{\rho\sigma}^M \Gamma_{\nu\lambda}^\sigma u^\nu u^\lambda v^\rho \end{aligned}$$

$$\rightarrow \frac{\partial}{\partial z} u^M = - \Gamma_{\nu\lambda}^M u^\nu u^\lambda$$

~~$$\frac{\partial}{\partial z} \left(\frac{\partial u^M}{\partial h} \right) = - \frac{\partial}{\partial z} \left(\Gamma_{\nu\lambda}^M u^\nu u^\lambda \right) = - \left(\frac{\partial \Gamma_{\nu\lambda}^M}{\partial z} v^\sigma u^\nu u^\lambda + \Gamma_{\nu\lambda}^M \frac{\partial u^\nu}{\partial z} u^\lambda + \Gamma_{\nu\lambda}^M u^\nu \frac{\partial u^\lambda}{\partial z} + \Gamma_{\nu\lambda}^M u^\nu u^\lambda \right)$$~~

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$$\begin{aligned} \frac{D^2 v^M}{D\tau^2} &= -\frac{\partial}{\partial h} (\Gamma_{\nu\lambda}^M u^\nu u^\lambda) + \partial_\sigma \Gamma_{\nu\lambda}^M u^\sigma u^\nu u^\lambda \\ &+ \Gamma_{\nu\lambda}^M (-\Gamma_{\rho\sigma}^\nu u^\rho u^\sigma) v^\lambda + 2 \Gamma_{\nu\lambda}^M u^\nu \cancel{u^\lambda} \frac{\partial u^\lambda}{\partial h} \\ &+ \Gamma_{\rho\sigma}^M \Gamma_{\nu\lambda}^\sigma u^\rho u^\nu v^\lambda \end{aligned}$$

$$\begin{aligned} \frac{\partial}{\partial h} (\Gamma_{\nu\lambda}^M u^\nu u^\lambda) &= \partial_\sigma \Gamma_{\nu\lambda}^M v^\sigma u^\nu u^\lambda + \Gamma_{\nu\lambda}^M \frac{\partial}{\partial h} u^\nu u^\lambda + \Gamma_{\nu\lambda}^M u^\nu \frac{\partial u^\lambda}{\partial h} \\ &= \cancel{\partial_\sigma \Gamma_{\nu\lambda}^M v^\sigma u^\nu u^\lambda} + \cancel{\Gamma_{\nu\lambda}^M v^\sigma} \end{aligned} \quad \text{cancel.}$$

$$\begin{aligned} \Rightarrow \frac{D^2 v^M}{D\tau^2} &= (\partial_\rho \Gamma_{\nu\lambda}^M - \partial_\lambda \Gamma_{\nu\rho}^M) u^\rho u^\nu v^\lambda \\ &+ (\Gamma_{\rho\sigma}^M \Gamma_{\nu\lambda}^\sigma - \Gamma_{\sigma\lambda}^M \Gamma_{\rho\nu}^\sigma) u^\rho u^\nu v^\lambda \\ &= R^M_{\nu\rho\lambda} u^\nu u^\rho v^\lambda \end{aligned}$$

2nd order derivative of v^M is proportional to the Riemann curvature tensor.



The separation between the geodesics changes.

v^M can be viewed as the relative coordinate between two close geodesic

$$\Delta x^M(\tau, h+oh) - x^M(\tau, h) = v^M.$$

③

Properties of the curvature tensor

• $[D_\mu D_\nu] A^\rho = R^\rho_{\ \sigma\mu\nu} A^\sigma$, hence $R^\rho_{\ \sigma\mu\nu} = -R^\rho_{\ \sigma\nu\mu}$

• Consider $\phi = \xi_\mu \xi^\mu = g_{\mu\nu} \xi^\mu \xi^\nu$

$$D_\lambda D_\rho \phi = D_\lambda (2 g_{\mu\nu} (D_\rho \xi^\mu) \xi^\nu) = 2 g_{\mu\nu} (D_\lambda D_\rho \xi^\mu) \xi^\nu + 2 g_{\mu\nu} D_\rho \xi^\mu D_\lambda \xi^\nu$$

$$\Rightarrow [D_\lambda, D_\rho] \phi = 2 g_{\mu\nu} [D_\lambda, D_\rho] \xi^\mu \xi^\nu$$

$$= 2 g_{\mu\nu} \xi^\nu R^\mu_{\ \sigma\lambda\rho} \xi^\sigma = 2 \xi^\nu \xi^\sigma R_{\nu\sigma\lambda\rho}$$

Since ϕ is a scalar, $[D_\lambda, D_\rho] \phi = 0 \Rightarrow \xi^\nu \xi^\sigma R_{\nu\sigma\lambda\rho} = 0$

$$\Rightarrow R_{\nu\sigma\lambda\rho} = -R_{\sigma\nu\lambda\rho}$$

• We define

$$\xi_{[\mu;\nu;\lambda]} = [\xi_{\mu;\nu;\lambda} - \xi_{\mu;\lambda;\nu}]$$

$$+ [\xi_{\nu;\lambda;\mu} - \xi_{\nu;\mu;\lambda}] + [\xi_{\lambda;\mu;\nu} - \xi_{\lambda;\nu;\mu}]$$

$$= [D_\lambda, D_\nu] \xi_\mu + [D_\mu, D_\lambda] \xi_\nu + [D_\nu, D_\mu] \xi_\lambda$$

$$= -R^\sigma_{\ \mu\lambda\nu} \xi_\sigma - R^\sigma_{\ \nu\mu\lambda} \xi_\sigma - R^\sigma_{\ \lambda\nu\mu} \xi_\sigma$$

$$A_{[\mu\nu\lambda]} = D_\lambda A_{\mu\nu} + D_\mu A_{\nu\lambda} + D_\nu A_{\lambda\mu}$$

$$A_{\mu\nu} = D_\mu A_\nu - D_\nu A_\mu = \partial_\mu A_\nu - \Gamma_{\mu\nu}^\lambda A_\lambda - (\partial_\nu A_\mu - \Gamma_{\nu\mu}^\lambda A_\lambda)$$

$$= \partial_\mu A_\nu - \partial_\nu A_\mu \Rightarrow A_{\mu\nu} = -A_{\nu\mu}$$

$$\Rightarrow \cancel{A_{[\mu\nu\lambda]} = D_\lambda A_{\mu\nu}}$$

$$D_\lambda A_{\mu\nu} = \partial_\lambda A_{\mu\nu} - \Gamma_{\lambda\mu}^\sigma A_{\sigma\nu} - \Gamma_{\lambda\nu}^\sigma A_{\mu\sigma} \quad \lambda \rightarrow \mu, \mu \rightarrow \nu, \nu \rightarrow \lambda$$

$$D_\mu A_{\nu\lambda} = \partial_\mu A_{\nu\lambda} - \Gamma_{\mu\nu}^\sigma A_{\sigma\lambda} - \Gamma_{\mu\lambda}^\sigma A_{\nu\sigma}$$

$$D_\nu A_{\lambda\mu} = \cancel{\partial_\nu A_{\lambda\mu}} - \Gamma_{\nu\lambda}^\sigma A_{\sigma\mu} - \Gamma_{\nu\mu}^\sigma A_{\lambda\sigma}$$

$$\text{if } A_{\mu\nu} = -A_{\nu\mu} \Rightarrow D_\lambda A_{\mu\nu} + D_\mu A_{\nu\lambda} + D_\nu A_{\lambda\mu} = \partial_\lambda A_{\mu\nu} + \partial_\mu A_{\nu\lambda} + \partial_\nu A_{\lambda\mu}$$

$$\begin{aligned} \text{then } A_{[\mu\nu\lambda]} &= \partial_\lambda [\partial_\mu A_\nu - \partial_\nu A_\mu] + \partial_\mu [\partial_\nu A_\lambda - \partial_\lambda A_\nu] \\ &\quad + \partial_\nu [\partial_\lambda A_\mu - \partial_\mu A_\lambda] \\ &= 0 \end{aligned}$$

$$\begin{aligned} \text{Then } \xi_{[\mu;\nu;\lambda]} &= D_\lambda D_\nu \xi_\mu - D_\nu D_\lambda \xi_\mu + D_\mu D_\lambda \xi_\nu - D_\lambda D_\mu \xi_\nu \\ &\quad + D_\nu D_\mu \xi_\lambda - D_\mu D_\nu \xi_\lambda \\ &= D_\lambda (D_\nu \xi_\mu - D_\mu \xi_\nu) + D_\mu (D_\lambda \xi_\nu - D_\nu \xi_\lambda) + D_\nu (D_\mu \xi_\lambda - D_\lambda \xi_\mu) \\ &= \partial_\lambda (\partial_\nu \xi_\mu - \partial_\mu \xi_\nu) + \partial_\mu (\partial_\lambda \xi_\nu - \partial_\nu \xi_\lambda) + \partial_\nu (\partial_\mu \xi_\lambda - \partial_\lambda \xi_\mu) \\ &= \partial_\lambda (\partial_\nu \xi_\mu - \partial_\mu \xi_\nu) + \partial_\mu (\partial_\lambda \xi_\nu - \partial_\nu \xi_\lambda) + \partial_\nu (\partial_\mu \xi_\lambda - \partial_\lambda \xi_\mu) \\ &= 0 \end{aligned}$$

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$$\Rightarrow R^\sigma_{\mu\lambda\nu} + R^\sigma_{\lambda\nu\mu} + R^\sigma_{\nu\mu\lambda} = 0$$

$$\text{or } R_{\sigma\mu\lambda\nu} + R_{\sigma\lambda\nu\mu} + R_{\sigma\nu\mu\lambda} = 0$$

$$\text{and } R_{\rho\mu\sigma\lambda} + R_{\rho\lambda\sigma\mu} + R_{\rho\sigma\mu\lambda} = 0.$$

- Check $R_{\mu\nu\lambda\rho} - R_{\lambda\rho\mu\nu} = ?$

$$\begin{aligned} R_{\mu\nu\lambda\rho} - R_{\lambda\rho\mu\nu} &= -(R_{\mu\lambda\rho\nu} + R_{\mu\rho\nu\lambda}) + \cancel{(R_{\lambda\mu\nu\rho} + R_{\lambda\nu\rho\mu})} \\ &\quad (R_{\mu\rho\nu\lambda} + R_{\nu\rho\lambda\mu}) \\ &= -R_{\mu\lambda\rho\nu} + R_{\nu\rho\lambda\mu} \end{aligned}$$

$$= (R_{\rho\lambda\nu\mu} + R_{\nu\lambda\mu\rho}) - (R_{\nu\lambda\mu\rho} + R_{\nu\mu\rho\lambda})$$

$$= R_{\rho\lambda\nu\mu} - R_{\nu\mu\rho\lambda} = R_{\lambda\rho\mu\nu} - R_{\mu\nu\lambda\rho}$$

$$\Rightarrow R_{\mu\nu\lambda\rho} - R_{\lambda\rho\mu\nu} = 0 \quad \text{or } R_{\mu\nu\lambda\rho} = R_{\lambda\rho\mu\nu}.$$

- Summary, $\begin{cases} R_{\mu\nu\lambda\rho} = -R_{\mu\rho\lambda\nu} = -R_{\nu\lambda\rho\mu} = R_{\lambda\rho\mu\nu} \\ R_{\mu\nu\lambda\rho} + R_{\mu\lambda\rho\nu} + R_{\mu\rho\nu\lambda} = 0 \end{cases}$

$$\left. \begin{array}{l} \mu\nu - \text{antisymmetric } 6 \\ \rho\lambda - \text{antisymmetric } 6 \end{array} \right\} \begin{array}{l} \cancel{6 \times 6 = 36} \\ \mu\nu; \rho\lambda \text{ symmetrized} \end{array} \left. \begin{array}{l} \\ \\ \end{array} \right\} \begin{array}{l} H_6^2 = \frac{6 \times 7}{2} \\ = 21 \end{array}$$

$$\text{Bianchi identity} \rightarrow \binom{n}{4} \quad \text{For } n=4, = 1 \Rightarrow 21 - 1 = 20$$

$$\text{For the general case, \# of degrees of freedom } \frac{n^2(n^2-1)}{12}$$

Contraction: \rightarrow Ricci tensor

$$R_{\mu\lambda} = g^{\nu\rho} R_{\mu\nu\lambda\rho} = R_{\mu\nu\lambda}{}^{\nu} = R_{\mu}{}^{\nu}{}_{\lambda\nu}$$

$$R_{\mu\lambda} = g^{\nu\rho} R_{\nu\mu\rho\lambda} = R^{\rho}{}_{\mu\rho\lambda} = R^{\rho}{}_{\mu\rho\lambda}$$

$$= R^{\nu}{}_{\mu\nu\lambda} = -R^{\nu}{}_{\mu\lambda\nu} \quad \text{---} \quad R^{\nu}{}_{\mu\lambda\nu}$$

Ricci scalar curvature $R = g^{\mu\lambda} R_{\mu\lambda}$

$R_{\mu\nu\lambda\rho}$ \rightarrow The only possible contraction to a scalar

Now we define a 2nd rank symmetric tensor from the

Ricci tensor

$$G^{\mu\nu} = R^{\mu\nu} - \frac{1}{2} g^{\mu\nu} R$$

Why define $G^{\mu\nu}$ in this form? $D_{\mu} G^{\mu\nu} = 0$,

which means that it is conserved!

Let's check: $[D_{\lambda} D_{\rho}] \xi^{\mu} = R^{\mu}{}_{\sigma\lambda\rho} \xi^{\sigma}$

$$\rightarrow [D_{\lambda}, D_{\rho}] \xi_{\mu} = R_{\mu\nu\lambda\rho} \xi^{\nu}$$

$$D_{\sigma} ([D_{\lambda}, D_{\rho}] \xi_{\mu}) = D_{\sigma} (R_{\mu\nu\lambda\rho} \xi^{\nu})$$

$$= D_{\sigma} (R_{\mu\nu\lambda\rho}) \xi^{\nu} + R_{\mu\nu\lambda\rho} D_{\sigma} \xi^{\nu} \quad (*)$$

on the other hand

$$[D_\lambda, D_\rho] D_\sigma \xi_\mu = R_{\sigma\nu\lambda\rho} D^\nu \xi_\mu + R_{\mu\nu\lambda\rho} D_\sigma \xi^\nu \quad (**)$$

Anti-symm (*) \rightarrow $(\sigma\lambda\rho)$ cyclic permutation

$$D_\sigma ([D_\lambda, D_\rho] \xi_\mu) + D_\lambda ([D_\rho, D_\sigma] \xi_\mu) + (D_\rho [D_\sigma, D_\lambda] \xi_\mu)$$

$$= D_\sigma (R_{\mu\nu\lambda\rho}) \xi^\nu + R_{\mu\nu\lambda\rho} D_\sigma \xi^\nu \quad \begin{matrix} \sigma \rightarrow \lambda \\ \lambda \rightarrow \rho \\ \rho \rightarrow \sigma \end{matrix}$$

$$+ D_\lambda (R_{\mu\nu\rho\sigma}) \xi^\nu + R_{\mu\nu\rho\sigma} D_\lambda \xi^\nu$$

$$+ D_\rho (R_{\mu\nu\sigma\lambda}) \xi^\nu + R_{\mu\nu\sigma\lambda} D_\rho \xi^\nu$$

(**) \rightarrow $(\sigma\lambda\rho)$ cyclic

$$[D_\lambda, D_\rho] D_\sigma \xi_\mu + [D_\rho, D_\sigma] D_\lambda \xi_\mu + [D_\sigma, D_\lambda] D_\rho \xi_\mu$$

$$= \cancel{R_{\lambda\nu\rho\sigma} D^\nu \xi_\mu} + R_{\sigma\nu\lambda\rho} D^\nu \xi_\mu + R_{\mu\nu\lambda\rho} D_\sigma \xi^\nu \\ + R_{\lambda\nu\rho\sigma} D^\nu \xi_\mu + R_{\mu\nu\rho\sigma} D_\lambda \xi^\nu \\ + R_{\rho\nu\sigma\lambda} D^\nu \xi_\mu + R_{\mu\nu\sigma\lambda} D_\rho \xi^\nu \quad \cancel{= 0}$$

cyclic relation $R_{\sigma\nu\lambda\rho} + R_{\lambda\nu\rho\sigma} + R_{\rho\nu\sigma\lambda} = 0$

$$\cancel{R_{\mu\nu\lambda\rho} + R_{\mu\nu\rho\sigma} + R_{\mu\nu\sigma\lambda} = 0}$$

\Rightarrow ~~on the other hand~~

$$\cancel{[D_\lambda, D_\rho] D_\sigma + [D_\rho, D_\sigma] D_\lambda + [D_\sigma, D_\lambda] D_\rho}$$

$$= \cancel{D_\sigma [D_\lambda, D_\rho] + D_\lambda [D_\rho, D_\sigma] + D_\rho [D_\sigma, D_\lambda]}$$

$$\begin{aligned}
& [D_\lambda, D_\rho] D_\sigma \xi_\mu + [D_\rho, D_\sigma] D_\lambda \xi_\mu + [D_\sigma, D_\lambda] D_\rho \xi_\mu \\
& = R_{\mu\nu\lambda\rho} D_\sigma \xi^\nu + R_{\mu\nu\rho\sigma} D_\lambda \xi^\nu + R_{\mu\nu\sigma\lambda} D_\rho \xi^\nu
\end{aligned}$$

Furthermore $[D_\lambda, D_\rho] D_\sigma + [D_\rho, D_\sigma] D_\lambda + [D_\sigma, D_\lambda] D_\rho$

$$= D_\sigma [D_\lambda, D_\rho] + D_\lambda [D_\rho, D_\sigma] + D_\rho [D_\sigma, D_\lambda]$$

$$\Rightarrow [D_\sigma R_{\mu\nu\lambda\rho} + D_\lambda R_{\mu\nu\rho\sigma} + D_\rho R_{\mu\nu\sigma\lambda}] \xi^\nu = 0$$

i.e. $D_\sigma R_{\mu\nu\lambda\rho} + D_\lambda R_{\mu\nu\rho\sigma} + D_\rho R_{\mu\nu\sigma\lambda} = 0$

$$g^{\nu\rho} [D_\sigma R_{\mu\nu\lambda\rho} + D_\lambda R_{\mu\nu\rho\sigma} + D_\rho R_{\mu\nu\sigma\lambda}] = 0$$

$$D_\sigma R_{\mu\lambda}{}^\sigma - D_\lambda R_{\mu\sigma}{}^\sigma - D_\rho R^{\rho\mu\sigma\lambda} = 0$$

$$g^{\mu\lambda} (D_\sigma R_{\mu\lambda}{}^\sigma - D_\lambda R_{\mu\sigma}{}^\sigma - D_\rho R^{\rho\mu\sigma\lambda}) = 0$$

$$D_\sigma R - D_\lambda R^\lambda{}_\sigma - D_\rho R^\rho{}_\sigma = 0$$

$$\frac{1}{2} D_\sigma R - D_\mu R^\mu{}_\sigma = 0$$

$$D_\mu R^\mu{}_\nu - \frac{1}{2} D_\nu R = 0$$

$$\Rightarrow D_\mu R^{\mu\nu} - \frac{1}{2} D^\nu R = 0$$

$$D_\mu (R^{\mu\nu} - \frac{1}{2} g^{\mu\nu} R) = 0 \Rightarrow D_\mu G^{\mu\nu} = 0$$

⊛ 2D curved space

$$R_{1212} = R_{2121} = -R_{1221} = -R_{2112}$$

$$g = g_{11}g_{22} - g_{12}g_{21} = \det g_{\mu\nu}$$

After complicated ~~by~~ but straightforward calculation, we have

$$R_{1212} = gK, \text{ where } K \text{ is the Gaussian curvature.}$$

$$\text{Then } R_{\mu\nu\lambda\rho} = K(g_{\mu\lambda}g_{\nu\rho} - g_{\mu\rho}g_{\nu\lambda})$$

$$\begin{aligned} \text{Ricci tensor } R_{\mu\nu} &= g^{\sigma\rho} R_{\mu\sigma\nu\rho} = g^{\sigma\rho} K(g_{\mu\nu}g_{\sigma\rho} - g_{\mu\rho}g_{\sigma\nu}) \\ &= K(2g_{\mu\nu} - g_{\mu\nu}) = Kg_{\mu\nu} \end{aligned}$$

$$R = g^{\mu\nu} R_{\mu\nu} = 2K$$

$$\frac{D^2 v^M}{D\tau^2} = R^M{}_{\nu\rho\lambda} u^\nu u^\rho v^\lambda$$

$$u^M = \frac{\partial x^M}{\partial \tau} \quad v^M = \frac{\partial x^M}{\partial h}$$

$$D^2 \delta x^M / D\tau^2 = R^M{}_{\nu\rho\lambda} \frac{dx^\nu}{d\tau} \frac{dx^\rho}{d\tau} \delta x^\lambda = K^M{}_{\lambda} \delta x^\lambda$$

$$\frac{d^2 x^i}{dt^2} = -\partial^i \phi|_x \rightarrow \boxed{\nabla^2 \phi(x) = 4\pi G_N \rho(x)}$$

$$\delta x^i = \tilde{x}^i - x^i$$

$$\frac{d^2 \tilde{x}^i}{dt^2} = -\partial^i \phi|_{\tilde{x}}$$

$$\frac{d^2 (\tilde{x}^i - x^i)}{dt^2} \simeq -\delta x^j \partial^i \partial_j \phi(x)|_x$$

$$\Rightarrow K^i{}_j = -\partial^i \partial_j \phi$$

$$\text{In empty} \quad \nabla^2 \phi = -\partial^i \partial_i \phi = K^i{}_i = 0$$

$$K^M{}_\mu = R^M{}_{\nu\rho\mu} \frac{dx^\nu}{d\tau} \frac{dx^\rho}{d\tau} = 0$$

$$\rightarrow -R_{\nu\rho} \frac{dx^\nu}{d\tau} \frac{dx^\rho}{d\tau} = 0 \Rightarrow \boxed{R_{\nu\rho} = 0} \rightarrow \text{free of matter}$$

in a space

\Rightarrow gravitational field equation free of matter

$$\boxed{R_{\mu\nu} = 0} \rightarrow \text{Curved space free of matter}$$

$$\dots \dots \dots \rho^M \dots \dots \dots = 0$$

weak field limit $g_{00}(x) = -(1 + 2\phi(x))$

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$$\Rightarrow \nabla^2 \phi = 4\pi G_N \rho \Rightarrow \nabla^2 g_{00} = 8\pi G_N T_{00} \rightarrow \begin{array}{l} \text{energy} \\ \text{- momentum} \\ \text{tensor} \end{array}$$

Conjecture
 $\rightarrow R^{\mu\nu} = \alpha T^{\mu\nu}$

Conservation law $D_\mu T^{\mu\nu} = 0$, but $D_\mu R^{\mu\nu} \neq 0$

Now $G^{\mu\nu} = R^{\mu\nu} - \frac{1}{2} g^{\mu\nu} R = \alpha T^{\mu\nu}$

$$D_\mu G^{\mu\nu} = \alpha D_\mu T^{\mu\nu} = 0$$

~~free~~ empty matter $\Rightarrow G^{\mu\nu} = 0 = R^{\mu\nu} - \frac{1}{2} g^{\mu\nu} R$

$$g_{\mu\nu} G^{\mu\nu} = R - \frac{4}{2} R = 0 \Rightarrow R = 0 \Rightarrow R^{\mu\nu} = 0$$

$$g_{\mu\nu} G^{\mu\nu} = \alpha g_{\mu\nu} T^{\mu\nu} = g_{\mu\nu} (R^{\mu\nu} - \frac{1}{2} g^{\mu\nu} R)$$

$$R - \frac{4}{2} R = \alpha T \Rightarrow R = -\alpha T$$

$$\Rightarrow \boxed{\begin{array}{l} R^{\mu\nu} = \alpha (T^{\mu\nu} - \frac{1}{2} g^{\mu\nu} T) \\ T = T^\mu{}_\mu \end{array}}$$

Action for Einstein's equation

The invariant volume in a curved manifold is $\int d^4x \sqrt{-g}$
 We need a scalar Lagrangian density based on metric.

If up to the order of quadratic terms, the only choice is R itself. Then up to a constant, we have

$$S_E = \beta \int d^4x \sqrt{-g} g^{\mu\nu} R_{\mu\nu} .$$

We consider the variation of metric under,

$$g_{\mu\nu} \rightarrow g_{\mu\nu} + \delta g_{\mu\nu}$$

$$\lim_{x \rightarrow \infty} g_{\mu\nu} \rightarrow \eta_{\mu\nu}$$

$$\lim_{x \rightarrow \infty} \delta g_{\mu\nu} \rightarrow 0$$

such that $\delta S_E = 0$.

$$g^{\mu\nu} g_{\nu\lambda} = \delta^\mu_\lambda \Rightarrow \delta g^{\mu\nu} g_{\nu\lambda} + g^{\mu\nu} \delta g_{\nu\lambda} = 0$$

$$\delta g^{\mu\nu} = - g^{\mu\sigma} \delta g_{\sigma\lambda} g^{\lambda\nu}$$

$$\begin{aligned} \delta(\sqrt{-g}) &= \frac{1}{2\sqrt{-g}} \delta(-g) = \frac{-1}{2\sqrt{-g}} \delta e^{\text{tr} \ln g} \\ &= \frac{-1}{2\sqrt{-g}} \cdot g \text{tr} [g^{-1} \delta g] = \frac{-g}{2\sqrt{-g}} g^{\mu\nu} \delta g_{\nu\mu} \end{aligned} \left. \begin{array}{l} = -\frac{\sqrt{-g}}{2} g_{\mu\nu} \delta g^{\nu\mu} \\ = \frac{\sqrt{-g}}{2} g^{\mu\nu} \delta g_{\nu\mu} \end{array} \right\}$$

$$\Gamma_{\mu\nu}^{\rho} = \frac{1}{2} g^{\rho\sigma} (\partial_{\mu} g_{\nu\sigma} + \partial_{\nu} g_{\sigma\mu} - \partial_{\sigma} g_{\mu\nu})$$

$$\delta \Gamma_{\mu\nu}^{\rho} = \frac{1}{2} \delta g^{\rho\sigma} (\partial_{\mu} g_{\nu\sigma} + \partial_{\nu} g_{\sigma\mu} - \partial_{\sigma} g_{\mu\nu}) + \frac{1}{2} g^{\rho\sigma} (\partial_{\mu} \delta g_{\nu\sigma} + \partial_{\nu} \delta g_{\sigma\mu} - \partial_{\sigma} \delta g_{\mu\nu})$$

The first line $-\frac{1}{2} g^{\rho\alpha} \delta g_{\alpha\beta} g^{\beta\sigma} (\partial_{\mu} g_{\nu\sigma} + \partial_{\nu} g_{\sigma\mu} - \partial_{\sigma} g_{\mu\nu})$

$$= -g^{\rho\alpha} \delta g_{\alpha\beta} \Gamma_{\mu\nu}^{\beta}$$

The 2nd line

$$D_{\mu} \delta g_{\nu\sigma} = \partial_{\mu} \delta g_{\nu\sigma} - \Gamma_{\mu\nu}^{\beta} \delta g_{\beta\sigma} - \Gamma_{\mu\sigma}^{\beta} \delta g_{\nu\beta}$$

$$D_{\nu} \delta g_{\sigma\mu} = \partial_{\nu} \delta g_{\sigma\mu} - \Gamma_{\nu\sigma}^{\beta} \delta g_{\beta\mu} - \Gamma_{\nu\mu}^{\beta} \delta g_{\sigma\beta}$$

$$D_{\sigma} \delta g_{\mu\nu} = \partial_{\sigma} \delta g_{\mu\nu} - \Gamma_{\sigma\mu}^{\beta} \delta g_{\beta\nu} - \Gamma_{\sigma\nu}^{\beta} \delta g_{\mu\beta}$$

$$\Rightarrow D_{\mu} \delta g_{\nu\sigma} + D_{\nu} \delta g_{\sigma\mu} - D_{\sigma} \delta g_{\mu\nu} = (\partial_{\mu} \delta g_{\nu\sigma} + \partial_{\nu} \delta g_{\sigma\mu} - \partial_{\sigma} \delta g_{\mu\nu}) - 2 \Gamma_{\mu\nu}^{\beta} \delta g_{\beta\sigma}$$

Then $\frac{1}{2} g^{\rho\sigma} (\partial_{\mu} \delta g_{\nu\sigma} + \partial_{\nu} \delta g_{\sigma\mu} - \partial_{\sigma} \delta g_{\mu\nu})$

$$= \frac{1}{2} g^{\rho\sigma} (D_{\mu} \delta g_{\nu\sigma} + D_{\nu} \delta g_{\sigma\mu} - D_{\sigma} \delta g_{\mu\nu}) + g^{\rho\sigma} \Gamma_{\mu\nu}^{\beta} \delta g_{\beta\sigma}$$

$$\Rightarrow \delta \Gamma_{\mu\nu}^{\rho} = \frac{1}{2} g^{\rho\sigma} (D_{\mu} \delta g_{\nu\sigma} + D_{\nu} \delta g_{\sigma\mu} - D_{\sigma} \delta g_{\mu\nu})$$

hence although $\Gamma_{\mu\nu}^{\rho}$ is not a tensor $\therefore \delta \Gamma_{\mu\nu}^{\rho}$ is a tensor.

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$$R^M{}_{\nu\lambda\rho} = \partial_\lambda \Gamma_{\nu\rho}^M - \partial_\rho \Gamma_{\nu\lambda}^M + \Gamma_{\lambda\sigma}^M \Gamma_{\nu\rho}^\sigma - \Gamma_{\rho\sigma}^M \Gamma_{\nu\lambda}^\sigma$$

$$R_{\nu\rho} = R^M{}_{\nu\mu\rho} = \partial_\mu \Gamma_{\nu\rho}^M - \partial_\rho \Gamma_{\nu\mu}^M + \Gamma_{\mu\sigma}^M \Gamma_{\nu\rho}^\sigma - \Gamma_{\rho\sigma}^M \Gamma_{\nu\mu}^\sigma$$

$$\delta R_{\nu\rho} = \partial_\mu \delta \Gamma_{\nu\rho}^M - \partial_\rho \delta \Gamma_{\nu\mu}^M + \delta \Gamma_{\mu\sigma}^M \Gamma_{\nu\rho}^\sigma + \Gamma_{\lambda\sigma}^M \delta \Gamma_{\nu\rho}^\sigma - \delta \Gamma_{\rho\sigma}^M \Gamma_{\nu\mu}^\sigma - \Gamma_{\rho\sigma}^M \delta \Gamma_{\nu\mu}^\sigma$$

$$= (\partial_\mu \delta \Gamma_{\nu\rho}^M + \Gamma_{\mu\sigma}^M \delta \Gamma_{\nu\rho}^\sigma - \Gamma_{\mu\nu}^\sigma \delta \Gamma_{\sigma\rho}^M - \Gamma_{\mu\rho}^\sigma \delta \Gamma_{\nu\sigma}^M)$$

$$- (\partial_\rho \delta \Gamma_{\nu\mu}^M + \Gamma_{\rho\sigma}^M \delta \Gamma_{\nu\mu}^\sigma - \Gamma_{\rho\nu}^\sigma \delta \Gamma_{\sigma\mu}^M - \Gamma_{\rho\mu}^\sigma \delta \Gamma_{\nu\sigma}^M)$$

$$= D_\mu (\delta \Gamma_{\nu\rho}^M) - D_\rho (\delta \Gamma_{\nu\mu}^M)$$

$$\Rightarrow \delta R_{\mu\nu} = D_\lambda \delta \Gamma_{\mu\nu}^\lambda$$

$$R = g^{\mu\nu} R_{\mu\nu}$$

$$- D_\nu \delta \Gamma_{\mu\lambda}^\lambda$$

$$\delta R = \delta g^{\mu\nu} R_{\mu\nu} + g^{\mu\nu} \delta R_{\mu\nu}$$

$$= -g^{\mu\lambda} g^{\rho\nu} \delta g_{\lambda\rho} R_{\mu\nu} + g^{\mu\nu} (D_\lambda \delta \Gamma_{\mu\nu}^\lambda - D_\nu \delta \Gamma_{\mu\lambda}^\lambda)$$

$$= -\delta g_{\lambda\rho} R^{\lambda\rho} + D_\lambda (g^{\mu\nu} \delta \Gamma_{\mu\nu}^\lambda) - D_\nu (g^{\mu\nu} \delta \Gamma_{\mu\lambda}^\lambda)$$

$$D_\lambda (g^{\mu\nu} \delta \Gamma_{\mu\nu}^\lambda) = \partial_\lambda (g^{\mu\nu} \delta \Gamma_{\mu\nu}^\lambda) + \Gamma_{\lambda\sigma}^\lambda (g^{\mu\nu} \delta \Gamma_{\mu\nu}^\sigma)$$

$$\Gamma_{\lambda\sigma}^\lambda = \frac{1}{2} g^{\lambda\tau} (\partial_\lambda g_{\sigma\tau} + \partial_\sigma g_{\tau\lambda} - \partial_\tau g_{\lambda\sigma})$$

$$\Gamma_{\lambda\sigma}^{\lambda} = \frac{1}{2} \overbrace{(g^{\tau} \partial_{\sigma} g) }^{\text{tr}} = \frac{1}{2} \partial_{\sigma} \overbrace{\text{tr} \ln g}^{\text{matrix}} = \frac{1}{2} \partial_{\sigma} \ln \det g \quad (4)$$

$$= \frac{1}{2} \partial_{\sigma} [\ln(-\det g)] = \dots \partial_{\sigma} \ln \sqrt{-\det g}$$

$$= \frac{1}{\sqrt{-\det g}} \partial_{\sigma} \sqrt{-\det g} \xrightarrow{\text{abbreviation}} \frac{1}{\sqrt{-g}} \partial_{\sigma} \sqrt{-g}$$

$$D_{\lambda} (g^{\mu\nu} \delta \Gamma_{\mu\nu}^{\lambda}) = \partial_{\lambda} (g^{\mu\nu} \delta \Gamma_{\mu\nu}^{\lambda}) + \frac{1}{\sqrt{-g}} \partial_{\lambda} (\sqrt{-g}) (g^{\mu\nu} \delta \Gamma_{\mu\nu}^{\lambda})$$

$$= \frac{1}{\sqrt{-g}} \partial_{\nu} (\sqrt{-g} \delta \Gamma_{\mu\nu}^{\lambda} g^{\mu\nu})$$

Similarly

$$\delta R = - \delta g_{\mu\nu} R^{\mu\nu} + \frac{1}{\sqrt{-g}} \partial_{\nu} (\sqrt{-g} g^{\mu\nu} \delta \Gamma_{\mu\nu}^{\lambda})$$

$$- \frac{1}{\sqrt{-g}} \partial_{\nu} (\sqrt{-g} g^{\mu\nu} \Gamma_{\mu\lambda}^{\lambda})$$

$$\Rightarrow \delta S_E = \beta \int d^4x (\delta \sqrt{-g}) R + \sqrt{-g} \delta R$$

$$= \beta \int d^4x \left(\frac{1}{2} \sqrt{-g} g^{\mu\nu} \delta g_{\mu\nu} R \right.$$

$$\left. + \sqrt{-g} (-\delta g_{\mu\nu} R^{\mu\nu} + \frac{1}{\sqrt{-g}} \partial_{\nu} (\sqrt{-g} (g^{\mu\nu} \delta \Gamma_{\mu\nu}^{\lambda} - g^{\mu\nu} \Gamma_{\mu\lambda}^{\lambda})) \right]$$

↙ total divergence

$$= \beta \int d^4x \sqrt{-g} \left(\frac{1}{2} g^{\mu\nu} R - R^{\mu\nu} \right) \delta g_{\mu\nu}$$

$$\delta S_E = 0 \Rightarrow R^{\mu\nu} - \frac{1}{2} g^{\mu\nu} R = G^{\mu\nu} = 0 \quad (5)$$

* In the presence of matter

$$S = S_E + S_{\text{matter}}$$

$$= \beta \int d^4x \sqrt{-g} R + \int d^4x \sqrt{-g} \mathcal{L}_{\text{matter}}$$

$$\delta S = \beta \int d^4x \sqrt{-g} \left(\frac{1}{2} g^{\mu\nu} R - R^{\mu\nu} \right) + \sqrt{-g} \left(\frac{1}{2} \delta g_{\mu\nu} T^{\mu\nu}_{\text{matter}} \right)$$

we assume $\delta S_{\text{matter}} = \frac{1}{2} \int d^4x \sqrt{-g} \delta g_{\mu\nu} T^{\mu\nu}$

then $\delta S = \int d^4x \sqrt{-g} \delta g_{\mu\nu} \left[\beta \left(\frac{1}{2} g^{\mu\nu} R - R^{\mu\nu} \right) + \frac{1}{2} T^{\mu\nu}_{\text{matter}} \right]$

$$\Rightarrow R^{\mu\nu} - \frac{1}{2} g^{\mu\nu} R = G^{\mu\nu} = \frac{1}{2\beta} T^{\mu\nu}_{\text{matter}}$$

$$G^{\mu\nu} = R^{\mu\nu} - \frac{1}{2} R g^{\mu\nu} = 8\pi G_N T^{\mu\nu}_{\text{matter}}$$

$$\beta = \frac{1}{16\pi G_N} \Rightarrow$$

If adding back 'c', we have.

$$S_E = \frac{c^3}{16\pi G_N} \int d^4x \sqrt{-g} R$$

$$\Rightarrow \frac{S_E}{\hbar} = \frac{l_p^{-2}}{16\pi} \int d^4x \sqrt{-g} R$$

Planck length

$$\Delta x \sim \frac{\hbar}{p} + \frac{GM}{c^2}$$

$$= \frac{\hbar c}{E} + \frac{GE}{c^4} \geq 2 \sqrt{\frac{\hbar G}{c^3}}$$

$$l_p^2 = \frac{\hbar G}{c^3} \Rightarrow \frac{c^3}{G} = \hbar l_p^{-2}$$

$$\text{or } \boxed{G = \frac{c^3}{\hbar} l_p^2}$$